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Considerese el movimiento de un fotón en el espacio-tiempo de Schwarzschild. Rezóñese que éste puede moverse según una órbita circular de radio r constante. Calúñese el radio de dicha órbita. Calúñese el periodo de la misma en el tiempo coordenado. Calúñese el periodo en el tiempo propio de un observador no geodráctico que se encuentra en un punto de la órbita.

Being massless, the photon's velocity satisfies

$$g_{\mu\nu} u^\mu u^\nu = 0$$

$$-f(r)\left(\frac{dt}{d\tau}\right)^2 + \frac{1}{f(r)}\left(\frac{dr}{d\tau}\right)^2 + r^2\left(\frac{d\theta}{d\tau}\right)^2 + r^2\sin^2\theta\left(\frac{d\phi}{d\tau}\right)^2 = 0.$$

Since θ may be taken without loss of generality as $\theta = \frac{\pi}{2}$

and

$$f \frac{dt}{d\tau} = E = \text{const.}$$

$$r^2 \dot{\phi} = L = \text{const.}$$

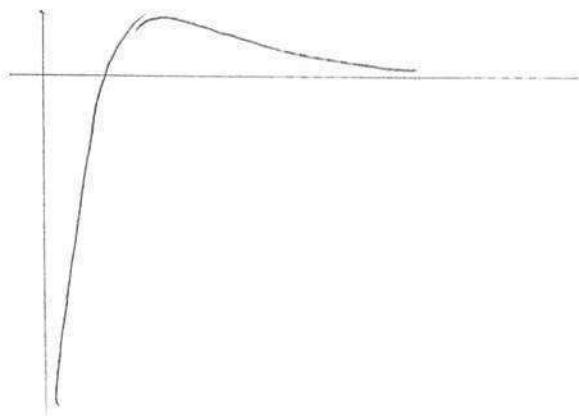
one has

$$-\frac{E^2}{f} + \frac{1}{f}\left(\frac{dr}{d\tau}\right)^2 + r^2 \frac{L^2}{r^4} = 0$$

$$\left(\frac{dr}{d\tau}\right)^2 + f \frac{L^2}{r^2} = E^2 \quad \Leftrightarrow \quad \left(\frac{dr}{d\tau}\right)^2 + \left(1 - \frac{2m}{r}\right) \frac{L^2}{r^2} = E^2 \quad \rightarrow$$

$$f = 1 - \frac{2m}{r}$$

$$V_{\text{eff}} = \frac{1}{r^2} \left(1 - \frac{2}{r} \right)$$



A circular orbit with constant r has $\frac{dr}{dt} = 0$. Since E^2 is constant for all r , the "effective potential energy"

$$V_{\text{eff}} := \left(1 - \frac{2m}{r}\right) \frac{L^2}{r^2}$$

must have a maximum. Let us find it

$$\begin{aligned} V'_{\text{eff}} &= \frac{2mL^2}{r^4} - \left(1 - \frac{2m}{r}\right) \frac{2L^2}{r^3} \\ V'_{\text{eff}}(r_0) &= 0 \end{aligned} \quad \left| \begin{array}{l} \frac{m}{r_0} = 1 - \frac{2m}{r_0} \Rightarrow r_0 = 3m \end{array} \right.$$

$$V''_{\text{eff}} = -\frac{24mL^2}{r^5} + \frac{6L^2}{r^4} = \frac{6L^2}{r^4} \left(1 - \frac{4m}{r}\right) \Rightarrow V''_{\text{eff}}(r_0) < 0 \Rightarrow$$

$\Rightarrow V_{\text{eff}}$ has a maximum at $r_0 = 3m$.

All in all

photon's circular orbit has radius $r_0 = 3m$

The coordinate period is

$$T_p = \int dt = \int_0^{2\pi} \frac{dt}{d\phi} d\phi = \int_0^{2\pi} \frac{E}{\frac{1-2m}{r_0} - \frac{L^2}{r_0^2}} d\phi = 27m^2 \frac{E}{L} 2\pi$$

On $r = r_0$ we have

$$T_p = T_p(r_0) \Rightarrow E^2 = f(r_0) \frac{L^2}{r_0^2} = \left(1 - \frac{2m}{r_0}\right) \frac{L^2}{r_0^2} \Rightarrow \frac{E}{L} = \frac{1}{\sqrt{3}r_0} = \frac{1}{3\sqrt{3}m}$$

Hence

$$T_p = 6\sqrt{3}\pi m$$

For an observer at a fixed point on the orbit, $\dot{r} = \dot{\phi} = 0$,
so that (recall that for the observer $u^2 = -1$)

$$-1 = -f(r_0) \left(\frac{dt}{d\tau} \right)^2$$

$$1 = \left(1 - \frac{2m}{r_0} \right) \left(\frac{dt}{d\tau} \right)^2 \Leftrightarrow d\tau = \frac{dt}{\sqrt{1 - \frac{2m}{r}}} \Rightarrow \tau_p = \frac{t_p}{\sqrt{3}}$$

$$\boxed{\tau_p = 6\pi m}$$